

Derivation of relaxation operators

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- 1 The Boltzmann equation
- 2 Method of moment relaxation
- 3 Shape of the set \mathcal{R}_m^+
- 4 ϕ divergence
- 5 Back to the model
- 6 Stationary shock wave

The Boltzmann equation

The Boltzmann equation

$$\partial_t f + v \cdot \nabla_x f = Q(f, f) = Q^+(f, f) - \nu(f)f, \quad (t, x, v) \in \mathbb{R}^+ \times \mathbb{R}^3 \times \mathbb{R}^3$$

Notations

$f(t, x, v)$: distribution function

$n = \int f \, dv$: number of molecules per unit volume

$nu = \int v f \, dv$: momentum ($m=1$)

$E = \frac{nu^2}{2} + \frac{3}{2}nT = \frac{1}{2} \int v^2 \, dv$: total energy

$\mathcal{H}(f) = \int (f \ln f - f) \, dv$: Boltzmann Entropy

1) $f(t, x, v) \geq 0$. Preservation of the positivity

2) Collision invariants (mass, momentum and energy)

$$\int Q(f, f)(1, v, v^2) dv = 0$$

3) $\exists \eta$ entropy density and

$$\mathcal{H}(f) = \int \eta(f) dv \quad \text{s.t.} \quad \int \eta'(f) Q(f, f) dv \leq 0.$$

3') For the Boltzmann equation $\eta(x) = x \ln(x) - x$,

$$\partial_t \mathcal{H}(f) + \operatorname{div} \int v \eta(f) dv \leq 0$$

4) Extended H theorem

$$\int_{\mathbb{R}^3} \eta'(f) Q(f, f) dv = 0 \Leftrightarrow Q(f, f) = 0 \Leftrightarrow \eta'(f) \in \text{Span}\{1, v, v^2\}$$

$$\Rightarrow f = \mathcal{M} = \frac{n}{(2\pi T)^{\frac{3}{2}}} \exp\left(-\frac{(v-u)^2}{2T}\right)$$

5) Correctness of the hydrodynamic limit

\Rightarrow Right properties on the linearized operator

Chapmann-Engskog expansion : $f = \mathcal{M}(1 + \varepsilon g) + O(\varepsilon^2)$

$$\partial_t f + v \cdot \nabla_x f = \frac{1}{\varepsilon} Q(f, f).$$

\Rightarrow Euler and Navier-Stokes (transport coefficients)

Summary

$f(t, x, v)$: distribution function

$$\underbrace{\frac{\partial f}{\partial t} + v \cdot \nabla_x f}_{\text{transport}} + \underbrace{F \cdot \nabla_v f}_{\text{Force term}} = \underbrace{C(f, f)}_{\text{Collision de term}}$$

n , u et T : density, velocity and temperature

$$n = \int_{\mathbb{R}^3} f dv, \quad u = \frac{1}{\rho} \int_{\mathbb{R}^3} v f dv, \quad T = \frac{1}{3\rho} \int_{\mathbb{R}^3} |v - u|^2 f dv.$$

Obtention of fluid models

- Equilibrium states : $C(f, f) = 0 \Leftrightarrow f = \mathcal{M}_f + O(\varepsilon^2)$
- $f = \mathcal{M}_f + O(\varepsilon)$ + moments extraction w.r.t. $(1, v, v^2) \Rightarrow$ Euler system
- $f = \mathcal{M}_f(1 + \varepsilon g) + O(\varepsilon^2)$ + moments extraction w.r.t. $(1, v, v^2) \Rightarrow$ Navier-Stokes system

Chapman-Enskog expansion

Parameter ε Knudsen number.

When $\varepsilon \rightarrow 0$, Boltzmann \Rightarrow fluid model

Rescaled Boltzmann equation

$$\partial_t f + v \cdot \nabla_x f = \frac{1}{\varepsilon} Q(f, f).$$

Chapman-Enskog expansion

- Equilibrium state : $Q(f, f) = 0 \Leftrightarrow f = \mathcal{M}_f$
- $f = \mathcal{M}_f$ + moments extraction w.r.t. $(1, v, v^2)$
 \Rightarrow Euler system
- $f = \mathcal{M}_f(1 + \varepsilon g)$ + moments extraction w.r.t. $(1, v, v^2)$
 \Rightarrow Navier-Stokes system

Euler system

Order 0

$$\left(\frac{\partial}{\partial t} + v \cdot \nabla_x\right) \mathcal{M} = \mathcal{M} \mathcal{L}_B(g) \quad (1)$$

with

$$\mathcal{L}_B(g) = \frac{1}{\mathcal{M}} Q(\mathcal{M}, \mathcal{M}g) + Q(\mathcal{M}g, \mathcal{M})$$

Integration of (1) w.r.t $(1, v, |v|^2) \Rightarrow$ Euler system

Euler system

$$\begin{aligned} \partial_t \rho + \operatorname{div}(\rho u) &= 0 \\ \partial_t(\rho u) + \operatorname{div}_x(\rho u \otimes u) + \nabla_x(\rho T) &= 0 \\ \partial_t\left(\rho\left(\frac{1}{2}|u|^2 + \frac{3}{2}T\right)\right) + \operatorname{div}_x\left(\rho u\left(\frac{1}{2}|u|^2 + \frac{5}{2}T\right)\right) &= 0. \end{aligned}$$

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Computation of g

Expression of **times derivatives** w.r.t **space derivatives**.

$$\left(\frac{\partial}{\partial t} + v \cdot \nabla_x\right) \mathcal{M} = (\mathbb{A}(V) : \mathbb{D}(u) - \mathbf{b}(V) \cdot \frac{\nabla_x T}{\sqrt{T}}) \mathcal{M} = \mathcal{M} \mathcal{L}_B(g)$$
$$V = \frac{v - u}{\sqrt{T}}.$$

Inversion of the relation $\Rightarrow g$

Sonine polynomials

$$\mathbb{A}(v) = v \otimes v - \frac{1}{3}|v|^2 Id, \quad \mathbf{b}(v) = \frac{v}{2}(v^2 - \frac{5}{2}).$$

$\mathbb{D}(u)$ (viscosity tensor) :

$$\mathbb{D}(u) = \frac{1}{2}(\nabla_x u + \nabla_x u^t) - \frac{1}{3} \operatorname{div}(u) Id.$$

Navier-Stokes system

Integration of $\left(\frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla_x\right)(\mathcal{M} + \varepsilon \mathcal{M}g)$ w.r.t $(1, \mathbf{v}, |\mathbf{v}|^2)$,

$$\partial_t \rho + \operatorname{div}_x(\rho \mathbf{u}) = 0$$

$$\partial_t(\rho \mathbf{u}) + \operatorname{div}_x(\rho \mathbf{u} \otimes \mathbf{u} + \rho T \operatorname{Id} - \varepsilon \mu \mathbb{D}(\mathbf{u})) = 0$$

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Transport Coefficients

$\mu = \mu(T, \rho, \mathbf{A}, \mathcal{L}_B^{-1})$: Viscosity, $\kappa = \kappa(T, \rho, \mathbf{b}, \mathcal{L}_B^{-1})$: Heat flux

Prandtl number

$$Pr = \frac{5}{2} \frac{\mu}{\kappa} \approx \frac{2}{3}.$$

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Collision operator

$$Q(f, f) = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} B(v - v_*, \omega) [f(t, x, v') f(t, x, v'_*) - f(t, x, v) f(t, x, v_*)] d\omega dv_*,$$

where

$$v' = v - \langle v - v_*, \omega \rangle \omega,$$

$$v'_* = v + \langle v - v_*, \omega \rangle \omega, \quad \omega \in \mathbb{S}^2$$

→ High complexity

⇒ Find a simplified operator

$$Q(f, f) \approx \lambda(G - f)$$

Method of moment relaxation

Relaxation operator

$$Q(f, f) \sim R(f) = \frac{1}{\tau}(\mathcal{M} - f), \quad \tau > 0$$

where \mathcal{M} is defined by

$$\mathcal{M}(v) = \frac{\rho}{(2\pi T)^{3/2}} \exp\left(-\frac{|v - u|^2}{2T}\right).$$

$$\mathcal{M} = \operatorname{Argmin}_{g \in C_f} \mathcal{H}(g)$$

where

$$C_f = \{g \geq 0 \text{ s.t. } \int_{\mathbb{R}^3} \begin{pmatrix} 1 \\ v \\ v^2 \end{pmatrix} g \, dv = \int_{\mathbb{R}^3} \begin{pmatrix} 1 \\ v \\ v^2 \end{pmatrix} f \, dv\}$$

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Properties of the BGK operator

Conservation laws

$$\int_{\mathbb{R}^3} (\mathcal{M} - f)(1, v, |v|^2) dv = (0, 0, 0),$$

Equilibrium states

$$\int_{\mathbb{R}^3} \rho(\mathcal{M} - f) \ln f dv = 0 \Leftrightarrow f = \mathcal{M},$$

H Theorem

$$\int_{\mathbb{R}^3} (\mathcal{M} - f) \ln f dv \leq 0.$$

Trend to equilibrium

$$\lim_{t \rightarrow +\infty} f(t) = \mathcal{M}.$$

Problem : Prandtl number not correct ≈ 1

Remark : Model coming from an entropy minimization problem

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Aim Construct relaxation operator

$$R(f) = \nu(G - f) \approx Q(f, f)$$

that is able to reproduce right the **transport coefficients**.

Requirements

- $R(f)$ « behaves as » linear operator
- "good properties" : Positivity, H-theorem, ...

Construction of the model

Linearization of f :

$$f = \mathcal{M}(1 + g) \implies Q(f, f) \approx \mathcal{M}\mathcal{L}_B(g)$$

\mathcal{L}_B : Linearized Boltzmann operator :

$$Q(\mathcal{M}(1 + g), \mathcal{M}(1 + g)) = \cancel{Q(\mathcal{M}, \mathcal{M})} + \mathcal{M}\mathcal{L}_B(g) + \cancel{Q(\mathcal{M}g, \mathcal{M}g)}$$

$\forall \phi$ test function,

$$\int_{\mathbb{R}^3} R(f)\phi dv \approx \int_{\mathbb{R}^3} \mathcal{M}\mathcal{L}_B(g)\phi dv = \int_{\mathbb{R}^3} \mathcal{M}g\mathcal{L}_B(\phi) dv = \int_{\mathbb{R}^3} f\mathcal{L}_B(\phi) dv$$

$$\Rightarrow \int_{\mathbb{R}^3} v(G - f)\phi dv = \int_{\mathbb{R}^3} f\mathcal{L}_B(\phi) dv$$

Relaxation constraints

Choice of a polynomial space in v : $\mathbb{P} = \text{span}(m_1(v), \dots, m_N(v))$,

s.t. $\mathbb{K} = \text{span}\{1, v, v^2\} \subset \mathbb{P}$

$(m_i)_i$ eigenvectors of \mathcal{L}_B : $\mathcal{L}_B(m_i) = -\nu_i m_i$

Test function : $\phi = m_i$

$$\int R(f) m_i(v) dv = \int \nu(G - f) m_i(v) dv = -\nu_i \int_{\mathbb{R}^3} f m_i(v) dv$$

$$\Rightarrow \int G m_i(v) dv = \left(1 - \frac{\nu_i}{\nu}\right) \int_{\mathbb{R}^3} f m_i(v) dv$$

$(-\nu_i)_{i=1,\dots,N}$ are nonpositive relaxation frequencies (eigenvalues)

First example

Set of tensors : $\mathbb{P} = \mathbb{K} \oplus^\perp \mathbb{A}$, where

$$\begin{aligned}\mathbb{K} &= \text{Span}\{1, v, v^2\} \\ \mathbb{A}(v) &= (v - u) \otimes (v - u) - \frac{1}{3} \|v - u\|^2 I_d\end{aligned}$$

Constraints on $G : C_f$

$$\begin{aligned}\int_{\mathbb{R}^3} G(1, v, v^2) dv &= \int_{\mathbb{R}^3} f(1, v, v^2) dv \\ \int_{\mathbb{R}^3} G \mathbb{A}(v - u) dv &= \left(1 - \frac{v \cdot \mathbb{A}}{v}\right) \int_{\mathbb{R}^3} f \mathbb{A}(v - u) dv\end{aligned}$$

Conservation laws $\Rightarrow v_i = 0$ on \mathbb{K} ($\mathbb{K} \subset \mathbb{P}$)

\Rightarrow Derivation of the ESBGK model [S.B., J. Schneider, 2008] with
 $\eta(f) = f \ln(f) - f$

Main example

Set of tensors : $\mathbb{P} = \mathbb{K} \oplus^\perp \mathbb{A} \oplus^\perp \mathbf{b}$, where

$$\mathbf{b}(v) = (v - u) \left(\frac{1}{2}(v - u)^2 - \frac{5}{2}T \right)$$

Constraints : C_f

$$\begin{aligned} \int_{\mathbb{R}^3} G(1, v, v^2) dv &= \int_{\mathbb{R}^3} f(1, v, v^2) dv \\ \int_{\mathbb{R}^3} G \mathbb{A}(v - u) dv &= \left(1 - \frac{\nu_{\mathbb{A}}}{\nu}\right) \int_{\mathbb{R}^3} f \mathbb{A}(v - u) dv \\ \int_{\mathbb{R}^3} G \mathbf{b}(v - u) dv &= \left(1 - \frac{\nu_{\mathbf{b}}}{\nu}\right) \int_{\mathbb{R}^3} f \mathbf{b}(v - u) dv \end{aligned}$$

Conservation laws $\Rightarrow \nu_i = 0$ on \mathbb{K}

How to define ν , $\nu_{\mathbb{A}}$ and $\nu_{\mathbf{b}}$, G ?

Computation of G

$$G = \operatorname{argmin}_{g \in C_f} \int_{\mathbb{R}^3} \eta(g) \, dv$$

(C_f set of moments constraints)

If $f \mapsto G$ is sufficiently smooth and if $R(f)$ is well-posed ("weak" H theorem)

$$\begin{aligned} \exists \eta, \quad \int \partial_f \eta(f) R(f) \, d\mathbf{v} &\leq 0, \\ R(f) = 0 &\Leftrightarrow \int \partial_f \eta(f) R(f) \, d\mathbf{v} = 0 \Leftrightarrow \partial_f \eta(f) \in \operatorname{span}\{\mathbf{1}, \mathbf{v}, \mathbf{v}^2\}. \end{aligned}$$

Chapman-Enskog expansion

$\Rightarrow O(\varepsilon)$: Euler and $O(\varepsilon^2)$: Navier-Stokes

Linearized operator/transport coefficients

Linearized operator of the Relaxation operator

$$\mathcal{L}_R(g) = \nu \left(\sum (1 - \frac{\nu_i}{\nu}) \mathbb{P}_{m_i} + \mathbb{P}_K - Id \right) (g)$$

Transport coefficients for the relaxation operator

$$\begin{aligned} \mu_R &= -\frac{k_B T}{10} \left\langle \mathcal{L}_R^{-1}(\mathbb{A}), \mathbb{A} \right\rangle = \frac{nk_B T}{\nu_A}, \\ \kappa_R &= -\frac{1}{3k_B T^2} \left\langle \mathcal{L}_R^{-1}(\mathbf{b}), \mathbf{b} \right\rangle = \frac{5}{2} \frac{nk_B^2 T}{m\nu_b}. \end{aligned}$$

$\langle \cdot, \cdot \rangle : L^2(\mathcal{M})$ dot product with the full contraction for tensor.

Aim : Recover the right viscosity and the heat conductivity

$$\mu_R = \mu_{\text{ref}} = \mu_B \text{ and } \kappa_R = \kappa_{\text{ref}} = \kappa_B$$

Definition of $\nu_{\mathbb{A}}$ and $\nu_{\mathbf{b}}$

Transport coefficients for Boltzmann

$$\mu_B = -\frac{k_B T}{10} \left\langle \mathcal{L}_B^{-1}(\mathbb{A}), \mathbb{A} \right\rangle, \quad \kappa_B = -\frac{1}{3k_B T^2} \left\langle \mathcal{L}_B^{-1}(\mathbf{b}), \mathbf{b} \right\rangle$$

Definition of $\nu_{\mathbb{A}}$ and $\nu_{\mathbf{b}}$

$$\mu_R = \mu_B, \quad \kappa_R = \kappa_B \Rightarrow \nu_{\mathbb{A}} = \frac{nT}{\mu_B}, \quad \nu_{\mathbf{b}} = \frac{5nT}{2\kappa_B} \Rightarrow Pr = \frac{5\mu}{2\kappa} = \frac{\nu_{\mathbf{b}}}{\nu_{\mathbb{A}}}$$

Remark : R is designed such that $\mathcal{L}_R^{-1} \sim \mathcal{L}_B^{-1}$ and **not** $\mathcal{L}_R \sim \mathcal{L}_B$

- ❶ What is the shape of the set of realizable moments

$$\mathcal{R}_{\mathbf{m}}^+ = \left\{ \int_{\mathbb{R}^3} f \mathbf{m}(v) dv, f \geq 0 \text{ a.e., } \int_{\mathbb{R}^3} f |m_i(v)| dv < +\infty \right\}$$

- ❷ Relaxation frequencies $\nu, (\nu_i)_i$ are such that $C_f \neq \emptyset$?

- ❸ Optimization problem : choose η such that

- Existence of a (unique) minimizer
- H theorem, positivity ...

No solution (in general) when $\eta(x) = x \ln(x)$ under the constraints

$$\begin{aligned}\int_{\mathbb{R}^3} g(1, \mathbf{v}, \mathbf{v}^2) dv &= \int_{\mathbb{R}^3} f(1, \mathbf{v}, \mathbf{v}^2) dv \\ \int_{\mathbb{R}^3} g \mathbb{A}(\mathbf{v} - \mathbf{u}) dv &= \left(1 - \frac{\lambda_{\mathbb{A}}}{\nu}\right) \int_{\mathbb{R}^3} f \mathbb{A}(\mathbf{v} - \mathbf{u}) dv \\ \int_{\mathbb{R}^3} g \mathbf{b}(\mathbf{v} - \mathbf{u}) dv &= \left(1 - \frac{\lambda_{\mathbf{b}}}{\nu}\right) \int_{\mathbb{R}^3} f \mathbf{b}(\mathbf{v} - \mathbf{u}) dv\end{aligned}$$

Artificial condition on $\int g|\mathbf{v}|^4 dv$?

The problem might not be well posed?

See [Junk, 1998, 2000], [J.Schneider, 2004], [Hauck et al, 2008], [Pavan, 2011]

Variational principle

Choose a strictly convex function η with domain in \mathbb{R}_+ and set

$$\mathcal{H}(g) = \int \eta(g) dv$$

For $\rho_f = \int f \mathbf{m}(v) dv$, define $L(\rho_f)$, with the relaxation constraints

$$(L(\rho_f))_i = \left(1 - \frac{v_i}{v}\right) \int_{\mathbb{R}^3} f m_i(v) dv$$

Problem

For $\rho_f \in \mathbb{R}^q$ ($q = \dim(\text{span}\{m_i\})$), find if possible a function G such that

- 1 $\int G \mathbf{m}(v) dv = L(\rho_f)$
- 2 $H(G) = \min_{\int g \mathbf{m}(v) dv = L(\rho_f)} H(g).$

Shape of the set \mathcal{R}_m^+

Assume $\mathbf{m}(\mathbf{v}) := (\mathbf{m}_0(\mathbf{v}), \dots, \mathbf{m}_k(\mathbf{v}), \dots, \mathbf{m}_n(\mathbf{v}))^T$: be a vector of tensors polynomial in $\mathbf{v} \in \mathbb{R}^3$ (pseudo Haar basis)

Realizability problem : $\rho = (\rho_0, \dots, \rho_n)$ list of tensor.

Is there a function $f \geq 0$ in $L^1(\mathbb{R}^d)$ s.t.

$$\int f m_i(\mathbf{v}) d\mathbf{v} = \rho_i \quad ?$$

Hamburger, Riesz, Haviland, Curto-Fialkow, Junk, Lasserre, Pichard, ...

Junk's theorem

Theorem (Junk, 2000)

- ① $\rho \in \mathcal{R}_m^+ \setminus \{0\} \Leftrightarrow \forall \alpha \neq 0 \text{ s.t. } \alpha \cdot m(v) \geq 0 \text{ a.e. there is } \rho \cdot \alpha > 0$
- ② $\mathcal{R}_m^{+,*}$ is an open convex set

Remark

$\mathcal{R}_m^+ \setminus \{0\}$ is characterized by the set of (negative) positive polynomials : all α such that $\alpha \cdot m(v) \geq 0$.

Issue : Characterize the set of nonnegative polynomials (tractable way ?)

Example

- ① "Euler" $\mathbf{m}(\mathbf{v}) = (1, \mathbf{v}, \mathbf{v}^2)$
- ② "Gauss" : $\mathbf{m}(\mathbf{v}) = (1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v})$
- ③ Grad : $\mathbf{m}(\mathbf{v}) = (1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v}, \mathbf{v}^2 \mathbf{v})$
- ④ Levermore : $\mathbf{m}(\mathbf{v}) = (1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v}, \mathbf{v}^2 \mathbf{v}, \mathbf{v}^4)$

XVIIth Hilbert's problem

Show that every nonnegative polynomial with coefficient in \mathbb{R} is a sum of square rational functions.

One of the important question about this problem :

If $p(\mathbf{v}) = \boldsymbol{\alpha} \cdot \mathbf{m}(\mathbf{v})$ is nonnegative, is it a sum of square (S.O.S) polynomials ?

Exemple : [Lasserre, 2009]

$$p(x) = 2x_1^4 + 2x_1^3x_2 - x_1^2x_2^2 + 5x_2^4 = \frac{1}{2}(2x_1^2 - 3x_2^2 + x_1x_2)^2 + \frac{1}{2}(x_2^2 + x_1x_2)^2$$

But $p(x) = x_1^2x_2^2(x_1^2 + x_2^2 - 1) + 1$ is nonnegative but not S.O.S.

Quadratic structured space

Example (Levermore space)

The Levermore space can be identified as a product of $(1, \mathbf{v}, \mathbf{v}^2) \vee (1, \mathbf{v}, \mathbf{v}^2) = (1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v}, \mathbf{v}^2 \mathbf{v}, \mathbf{v}^4)$

A square polynomial $P(v) = (a + \mathbf{b} \cdot \mathbf{v} + c\mathbf{v}^2)^2$ can be written as

$$\boldsymbol{\beta}^T M \boldsymbol{\beta}, \quad \text{with} \quad \boldsymbol{\beta} = (a, \mathbf{b}, c)^T \in \mathbb{R} \times \mathbb{R}^3 \times \mathbb{R}$$

$$M = \begin{bmatrix} 1 & \mathbf{v}^T & \mathbf{v}^2 \\ \mathbf{v} & \mathbf{v} \otimes \mathbf{v} & \mathbf{v}^2 \mathbf{v} \\ \mathbf{v}^2 & \mathbf{v}^2 \mathbf{v}^T & \mathbf{v}^4 \end{bmatrix}$$

What about Grad space ?

$(1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v}, \mathbf{v}^2 \mathbf{v})$ has no quadratic structure

Definition

For $f \geq 0$ in L^1_{Lev} define a Hankel matrix H as

$$H = \int_{\mathbb{R}^3} M f(v) dv.$$

$$\Rightarrow \int (a + \mathbf{b} \cdot \mathbf{v} + c\mathbf{v}^2)^2 f dv = \boldsymbol{\beta}^T H \boldsymbol{\beta} > 0$$

Necessary condition :

H must be definite positive.

Converse statement ?

True if every positive polynomial is a Sum Of Square (S.O.S.)

$$(\boldsymbol{\beta}^T H \boldsymbol{\beta} = \boldsymbol{\alpha} \cdot \boldsymbol{\rho})$$

Known results

Known results between positive polynomials and S.O.S in \mathbb{R}^d

- ① $d = 1$: every positive polynomial is a S.O.S
- ② $d = 2$: true for polynomial of degree $n \leq 4$ but not always if $n \geq 6$ (Hilbert 1893)
- ③ $d \geq 3$: true for polynomial of degree $n = 2$ but not always if $n \geq 4$

The first explicit counterexample for non S.O.S polynomial in dimension 2 was only found in 1966 !

Theorem

Artin (1927) Every nonnegative polynomial is a sum of square rational functions.

Grad space

Positive polynomials in **Grad space** $\in \text{span}(1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v})$ (Gauss space)

i.e. Every positive polynomial $\boldsymbol{\alpha} \cdot \mathbf{m}(\mathbf{v})$ writes as $(\boldsymbol{\beta}, 0) \cdot (1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v}, \mathbf{v}^2 \mathbf{v})$

In Gauss space, S.O.S \Leftrightarrow Non negative polynomials

Characterization by S.O.S. in the Gauss space **and** of realizable moment by the Hankel matrix

Proposition

$\boldsymbol{\rho} = (n, n\mathbf{u}, \Pi, Q) \in \mathcal{R}_{\text{Grad}}^+$ iff $n > 0, \Pi - n\mathbf{u} \otimes \mathbf{u} > 0$.

$$n\Pi = \int \mathbf{v} \otimes \mathbf{v} f \, d\mathbf{v}, \quad nQ = \int \mathbf{v}^2 \mathbf{v} f \, d\mathbf{v}.$$

Relaxation in Grad basis

$\text{span}(1, \mathbf{v}, \mathbf{v} \otimes \mathbf{v}, \mathbf{v}^2 \mathbf{v})$ is generated by

$$\mathbf{a}(\mathbf{v} - \mathbf{u}) := (1, (\mathbf{v} - \mathbf{u}), (\mathbf{v} - \mathbf{u})^2 - 3T, \mathbb{A}(\mathbf{v} - \mathbf{u}), \mathbf{b}(\mathbf{v} - \mathbf{u}))$$

Moments of f

$$\int_{\mathbb{R}^3} f \mathbf{a}(\mathbf{v} - \mathbf{u}) \, d\mathbf{v} = (n, 0, 0, \bar{\mathbb{P}}, \mathbf{q})$$

where $\bar{\mathbb{P}}$: traceless pressure tensor and \mathbf{q} : heat flux

Proposition

$$(n, 0, 0, \lambda_{\mathbb{A}} \bar{\mathbb{P}}, \lambda_{\mathbf{b}} \mathbf{q}) \in \mathcal{R}_{\text{Grad}}^{+,*} \quad \forall \lambda_{\mathbb{A}} \in [-\frac{1}{2}, 1] \text{ and } \forall \lambda_{\mathbf{b}} \in \mathbb{R}$$

Remark : The heat flux can take any value

Problem

Legendre dual function $h^* : \mathbb{R}^q \longrightarrow \overline{\mathbb{R}} = \mathbb{R} \cup \{+\infty, -\infty\} :$

$$\forall \alpha \in \mathbb{R}^q, \quad h^*(\alpha) = \sup_{\rho \in \mathbb{R}^q} (\rho \cdot \alpha - h(\rho)).$$

Example : $\phi(x) = x \ln(x) - x \Rightarrow \phi^*(y) = \exp(y)$

$$\int_{\mathbb{R}^3} (f \ln f - f) \, dv \implies \text{Maxwellian distribution function} : \mathcal{M}$$

Property : If η is convex, $(\eta')^{-1} = \eta^{*}$

Ex : If $\eta(x) = x \ln x - x$, $(\eta')^{-1} = \exp$, $\eta^* = \exp$, $(\eta^*)' = \exp$

Problem

Definition (Entropy, Entropy density)

$$\mathcal{H}(g) = \int \eta(g) dv, \quad h(\rho) = \min_{\int g \mathbf{m}(v) dv = \rho} \mathcal{H}(g).$$

Problem

For $\rho \in \text{dom}(h)$ find if possible a function G such that

- $\mathcal{H}(G) = h(\rho)$
- $\int G \mathbf{m}(v) dv = \rho$
- $G = (\eta')^{-1}(\alpha \cdot \mathbf{m}(v)) = (\eta^*)'(\alpha \cdot \mathbf{m}(v))$ (analytical form)
- $G = \mathcal{M}$ for $\rho = \int \mathcal{M} \mathbf{m}(v) dv$ (hydrodynamic limit)

ϕ divergence

« Renormalisation » map of [Abdel-Malik, Van Brummelen, 2015]

One starts from $(1 + \frac{x}{N})^N \rightarrow \exp(x)$ and looks for solutions of the form

$$G = \mathcal{M}(1 + \frac{\alpha \cdot m(v)}{N})_+^N,$$

$(x)_+ =$ positive part

Inverse function of $(1 + \frac{x}{N})_+^N : \widetilde{\ln}(y) = Ny^{1/N} - N \rightarrow \ln(y)$

\mathcal{H} is replaced by

$$\mathcal{H}_N = \int \mathcal{M} \phi_N(f/\mathcal{M}) dv, \quad \text{with } \phi_N(x) = x \widetilde{\ln}(x)$$

Remark : polynomial growth of $(1 + \frac{x}{N})^N$ instead of exponential

General result by convex analysis

Theorem

- ① $\phi : \mathbb{R} \rightarrow \mathbb{R} \cup \{+\infty\}$ *strictly convex and differentiable of \mathbb{R}^+*
- ② $\phi(0) = 0, \lim_{p \rightarrow 0} \frac{\phi(p)}{p} \in \mathbb{R}, \phi$ *superlinear*
- ③ $\forall \alpha, \phi^*(\alpha \cdot \mathbf{m}(v)) \in L^1(\mathcal{M} dv)$: **Csiszar assumption**

Then

$G = \mathcal{M}(\phi^*)'(\alpha \cdot \mathbf{m}(v))$: *unique solution to the variational problem :*

$$\mathcal{H}(G) = h(\rho) = \min_{\int g \mathbf{m}(v) dv = \rho} \mathcal{H}(g)$$

$\phi_N^*(x) = (1 + \frac{x}{N})_+^N$ satisfies Csiszar assumption

$\phi(x) = x \ln(x) - x$ with $\phi^*(x) = \exp(x)$ **does not** satisfy Csiszar assumption

Back to the model

In the Grad space

Step 1 : For $f \geq 0$, $f \in L^1_{Grad}$, consider $\rho_f = \int f \mathbf{a}(v - u) dv$

Step 2 : Relaxation :

$$v \geq v_A, v_b \geq 0$$

$$L(\rho_f) = (n, 0, 0, (1 - \frac{v_A}{v})\bar{P}, (1 - \frac{v_b}{v})\mathbf{q})$$

$$(n, 0, 0, \lambda_A \bar{P}, \lambda_b \mathbf{q}) \in \mathcal{R}_{Grad}^{+,*} \quad \forall \lambda_A \in [-\frac{1}{2}, 1] \text{ and } \forall \lambda_b \in \mathbb{R}$$

Step 3 : Solve the variational problem for some ϕ divergence

$$\exists ! \boldsymbol{\alpha}, \quad G = \mathcal{M}_f(\phi^*)'(\boldsymbol{\alpha} \cdot \mathbf{a}(v - u)),$$

with

$$\int G \mathbf{a}(v - u) dv = L(\rho_f), \quad \text{and} \quad \mathcal{H}(G) = h(L(\rho_f))$$

Equation

$$\partial_t f + \mathbf{v} \cdot \nabla_x f = \nu(G - f)$$

- 1 Positivity, conservation laws, H-theorem, Galilean invariance
- 2 Exact hydrodynamic limit if $\nu \geq \nu_A$, ν_b

$$\nu_A = \frac{nT}{\mu_B}, \quad \nu_b = \frac{5}{2} \frac{nT}{\kappa_B}$$

Remark

$\int f \ln f$ is almost a Lyapunov functional for the inhomogeneous equation if $\phi(x) \approx x \ln x$ (ex : $\phi_N(x) = N(x^{1+1/N} - x)$)

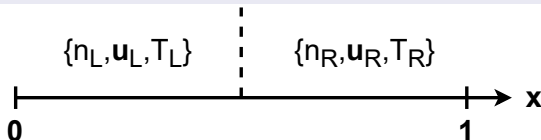
$$\int R(f) \phi'_N(f/M) dv \longrightarrow \int R(f) \ln(f/M) dv$$

Stationary shock wave

Stationary normal shock wave - 1

Purpose : the 1D domain is divided in 2 regions, with different gas states, and let the gas relaxed to the stationary state.

Domain initialization



Boundary conditions

- Left at 0 : $\{n_L, u_L, T_L\}$
- Right at 1 : $\{n_R, u_R, T_R\}$

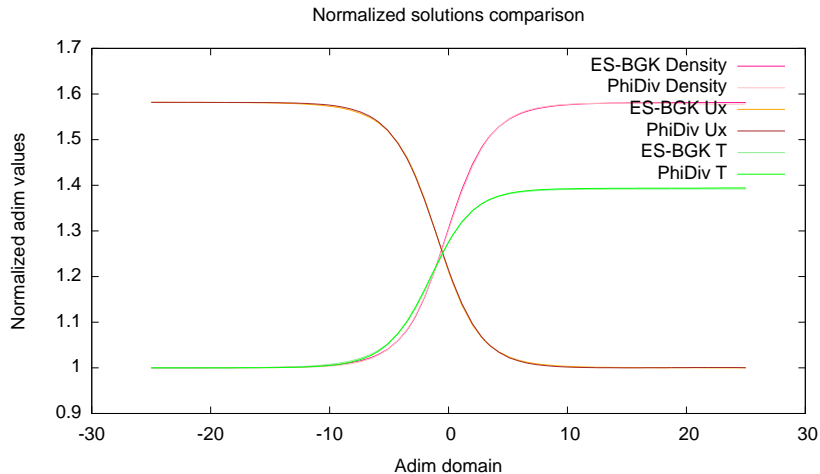
Stationary normal shock wave

Mach	Boundary	n	u	T
1.4	left	1	1.278	1
	right	1.581	0.808	1.392

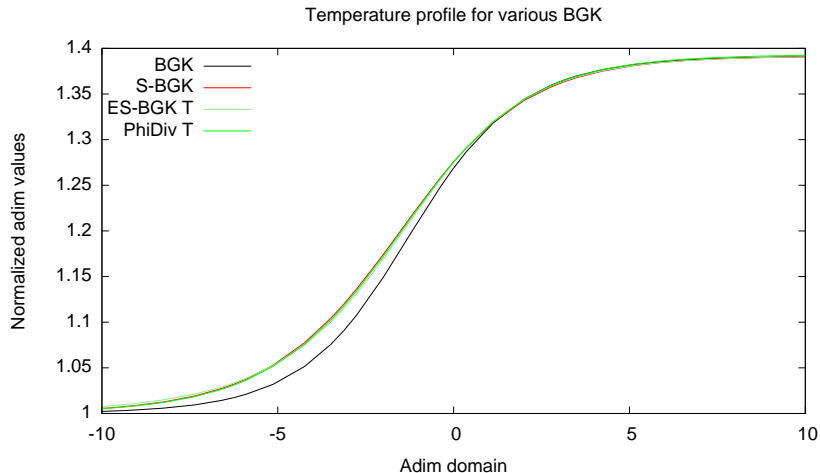
Code characteristics

- DVM on top of a Discontinuous Galerkin Advection solver
- 1D physical, 3D molecular velocities
- BGK model : BGK, S-BGK, ES-BGK, phi-div ($N = 3$)

Stationary normal shock wave - Results



Stationary normal shock wave - Results



- Fick matrix gas mixtures : [S.B, V. Pavan, J. Schneider., 2012]
- ESBGK models for mono and polyatomic gas mixtures :
viscosity (and shear viscosity), heat conductivity :
[S.B., J. Schneider, 2008] (mono), [S.B., J. Schneider, 2009] (poly),
[S.B., 2015, 2021]
- Polyatomic reacting gases, discrete energy, Fick matrix
[S.B., J.Schneider., 2014], [J.Schneider. 2015]
- Fick matrix poly (and mono) gas mixtures : 2 viscosities and Fick matrix
[S.B., Guillon, Thieullen, 2024]
- ESBGK model with a general formalism for microscopic energy
introduced by [Borsoni, Bisi, Groppi] : [S.B., Pollino, 2025]

- Relaxation model for multispecies (mono and poly) leading to the full set of transport coefficients (Phenomenological or Onsager matrix)
- Numerical method to simulate phi-divergence operators
Developpement of a code : [S. B., Y. Jobic, V. Pavan, J. Schneider]
In progress
- BGK model for Enskog.
[S. B., A. Derro, A. Takahashi]
In progress

THANKS FOR YOUR
ATTENTION!